

## MLS and CALIOP Cloud Ice Measurements in the Upper Troposphere: A Constraint from Microwave on Cloud Microphysics

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### ABSTRACT

This study examines the consistency and microphysics assumptions among satellite ice water content (IWC) retrievals in the upper troposphere with collocated A-Train radiances from Microwave Limb Sounder (MLS) and lidar backscatters from Cloud–Aerosol Lidar with Orthogonal Polarization (CALIOP). For the cases in which IWC values are small ( $<10 \text{ mg m}^{-3}$ ), the cloud ice retrievals are constrained by both MLS 240- and 640-GHz radiances and CALIOP 532-nm backscatter  $\beta_{532}$ . From the observed relationships between MLS cloud-induced radiance  $T_{\text{cir}}$  and the CALIOP backscatter integrated  $\gamma_{532}$  along the MLS line of sight, an empirical linear relation between cloud ice and the lidar backscatter is found:  $\text{IWC}/\beta_{532} = 0.58 \pm 0.11$ . This lidar cloud ice relation is required to satisfy the cloud ice emission signals simultaneously observed at microwave frequencies, in which ice permittivity is relatively well known. This empirical relationship also produces IWC values that agree well with the CALIOP, version 3.0, retrieval at values  $<10 \text{ mg m}^{-3}$ . Because the microphysics assumption is critical in satellite cloud ice retrievals, the agreement found in the  $\text{IWC}$ – $\beta_{532}$  relationships increase fidelity of the assumptions used by the lidar and microwave techniques for upper-tropospheric clouds.

### 1. Introduction

Cloud ice and occurrence frequency in the upper troposphere contribute significantly to Earth's total radiation and energy budgets. However, current climate and weather models produce a wide spread of values for these variables, leading to large uncertainties in the predicted dynamics and precipitation at the surface (e.g., Waliser et al. 2009; Eliasson et al. 2011; Jiang et al. 2012). Improving cloud ice retrieval and modeling is imperative and can be achieved by reducing uncertainty

of the assumptions about cloud microphysics in remote sensing and modeling physics (e.g., Su et al. 2013).

Primary sources of the observed ice cloud microphysical properties are in situ measurements from high-altitude field campaigns, but these data are limited to the types of clouds and systems accessible by aircrafts. As a result, further assumptions and extrapolation are needed for global cloud systems, so that a generalized parameterization can be used to retrieve and model cloud ice and other properties (e.g., Zhao and Weng 2002; Heymsfield et al. 2005; Delanoë and Hogan 2008; Austin et al. 2009). Because of large uncertainties associated with the assumption/parameterization on cloud microphysics, satellite cloud ice retrievals remain different by a factor of 2 or more (e.g., Wu et al. 2009).

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Multipatform multifrequency remote sensing now allows self-consistency evaluation of satellite cloud retrievals (e.g., cloud ice and cloud effective mass diameter  $D_{\text{mm}}$ ). For example, passive microwave radiometers at 89 and 150 GHz (Zhao and Weng 2002), multifrequency radars (Majurec 2008; Matrosov 2011), and joint lidar–radar systems (Delanoë and Hogan 2008) were among the attempts of this sort. Unlike satellite versus ground-based observations, the cross validation among satellite sensors can be made on a global basis, but this requires coincident and collocated measurements and an overlapped sensitivity between sensors. The coincident measurements are generally uncommon among spaceborne platforms, but the National Aeronautics and Space Administration’s (NASA) A-Train has collected an unprecedented amount of such measurements. NASA’s A-Train is a set of satellites with multiple sensors that fly in formation on a sun-synchronous orbit (e.g., L’Ecuyer and Jiang 2010). The lidar–radar approach for cloud ice retrievals is applicable to the A-Train Cloud–Aerosol Lidar with Orthogonal Polarization (CALIOP) and *CloudSat* measurements, because the collocated footprints are within 1-min separation in time, or 400 km in distance along the track. However, there is a narrow overlap between lidar and radar in sensitivity and penetration, limiting the dynamic range of joint ice cloud retrievals using *CloudSat* and CALIOP. The overlapped clouds are often seen as a narrow layer in the A-Train curtain plots, at the bottom (top) of CALIOP (*CloudSat*) cloud profiles. Beyond the overlapped sensitivity range, cloud ice retrievals depend heavily on the individual sensor, for example, CALIOP for tenure cirrus and *CloudSat* for cumulonimbus.

In this study we evaluate the consistency of upper-tropospheric ice water content (IWC) measurements from another pair of A-Train sensors: CALIOP 532-nm backscatter and Microwave Limb Sounder (MLS) 240- and 640-GHz radiances. These remote sensing techniques are based on two independent ice cloud signals, namely light scattering (CALIOP) and ice emission (MLS) methods. For optically thin cirrus the MLS sensitivity comes from the blackbody emission of cloud ice over a long limb path, whereas CALIOP relies on the backscattering of cloud ice particles. In the case where ice particles are too small ( $<10\ \mu\text{m}$ ) for *CloudSat* to detect cloud scattering signal ( $-30\ \text{dBZ}$ ), MLS can still detect the thermal emission of ice particles in the absence of scattering. The focus of our study is on the statistical consistency of satellite cloud ice measurements between a large ensemble of the datasets acquired at 15-km altitude after May 2008, when MLS and CALIOP samplings become nearly collocated and within 1 min apart along the A-Train curtain (Fig. 1).

## 2. Data and methods

*Aura* MLS, launched in July 2004, has seven radiometers with horizontal (H) and vertical (V) polarizations at frequencies near 118 (H, V), 190 (V), 240 (H), 640 (H) GHz and 2.5 (H, V) THz. MLS can detect ice clouds if the cloud top reaches MLS limb tangent heights  $h_t$ , but the sign of the cloud-induced radiance  $T_{\text{cir}}$ , the difference between the observed radiance and modeled clear-sky radiance after gas species retrievals are completed, is a function of  $h_t$  and cloud IWC or ice water path (IWP) in the limb line of sight (LOS; Wu et al. 2006, 2008, 2009). In the standard MLS cloud ice product, the 240-GHz  $T_{\text{cir}}$  is used for the IWC retrieval, while the  $T_{\text{cir}}$  at other frequencies are also computed and archived as diagnostic products. Comparing upper-tropospheric MLS 240-GHz and *CloudSat* IWC, Wu et al. (2009) found that MLS IWC retrieval is lower by as large as a factor of 5 at pressures  $<200\ \text{hPa}$ . This large difference was thought to be mainly because of different assumptions about ice microphysics between these retrievals. The MLS retrieval uses the parameterization formulated by McFarquhar and Heymsfield (1997, hereafter MH97), whereas *CloudSat* assumes a gamma size distribution in cloud ice retrieval (Austin et al. 2009). Another factor that can cause the lower IWC is MLS sensitivity limitation. MLS cloud signals can become saturated in a very large IWC case. As a result, it reports the large IWC for a smaller value, causing a low bias in the MLS retrieval as IWC increases, which is evident in the MLS–*CloudSat* probability density function (PDF) comparison (Wu et al. 2009).

CALIOP is a dual-wavelength (532 and 1064 nm) and dual-polarization (perpendicular and parallel at 532 nm) lidar (Winker et al. 2009) that has footprints collocated with *CloudSat* 94-GHz cloud profiling radar (CPR; Stephens et al. 2002). The original level-1  $\beta_{532}$  data have 583 vertical levels with resolutions from 30 m near the surface to 300 m in the stratosphere, and horizontal resolution of 300 m. CALIOP level-2 data [L2\_05 kmCPro, version 3.0 (V3.0), dataset] contain extinction and IWC profiles at 60-m vertical and 5-km horizontal resolutions, which are retrieved from the  $\beta_{532}$  measurements. A detailed discussion of the CALIOP extinction coefficient retrieval can be found in Young and Vaughan (2009), where the inversion from the backscatter profile takes into account two-way molecular/particulate transmittance. For cloud extinction and IWC retrievals, the CALIOP version 3.01 (V3.01) algorithm assumes  $25 \pm 10\ \text{sr}$  for the lidar ratio, 0.6 for multiple-scattering factor (Powell et al. 2010), and a temperature-invariant relation ( $\text{IWC} = 26\sigma^{1.22}$ ) for the extinction-to-IWC retrieval, where these coefficients correspond to the parameters at  $-73^\circ\text{C}$  in (Heymsfield et al. 2005). Because the extinction

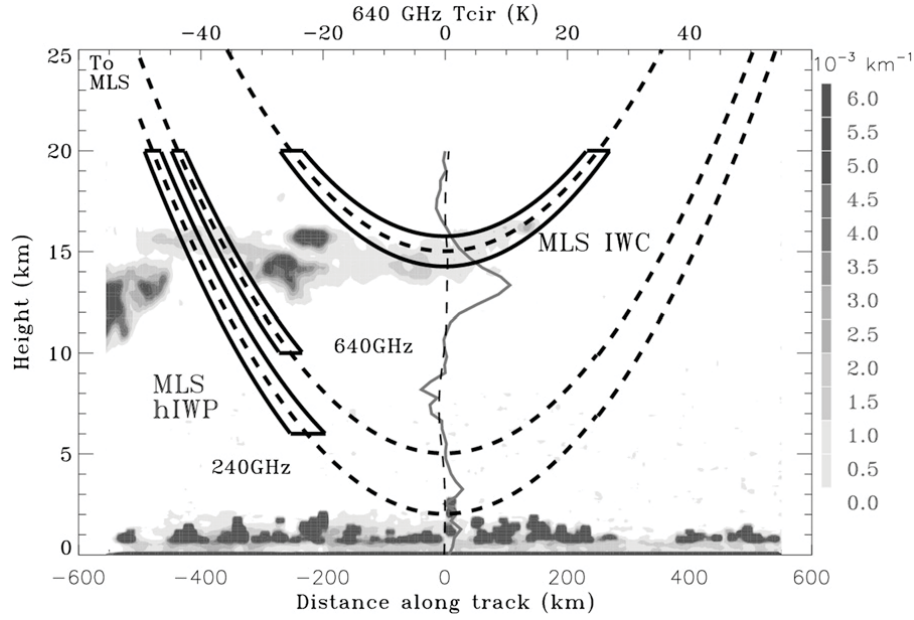


FIG. 1. A-Train MLS, CALIOP, and *CloudSat* cloud measurements on 19 May 2008 near 57.1°W, 13.5°N. The shaded contours are the attenuated CALIOP 532-nm backscatters. The MLS measurement volumes are indicated by the curved boxes. At high  $h_t$ , MLS radiances can penetrate through the atmosphere limb to measure IWC. At low  $h_t$ , MLS radiances cannot see through the tangent point but can be used to measure IWP. The observed MLS 640-GHz  $\Delta T_{\text{cir}}$  profile (the gray curve in the middle) shows a peak of approximately 10 K at  $h_t \sim 14$  km, whereas the dashed profile is the calculated 640-GHz  $\Delta T_{\text{cir}}$  from *CloudSat* IWC because the radar reports little IWC in this region because of the lack of sensitivity to small ice crystals.

and backscatter are linearly related, the CALIOP IWC is proportional to  $\beta_{532}^{1.22}$ . As shown in Fig. 2, for small ( $< 50 \times 10^{-3} \text{ km}^{-1} \text{ sr}^{-1}$ )  $\beta_{532}$  values, the retrieved IWC can be approximated by  $0.29\beta_{532}^{1.22}$ , where the linear slopes of 0.4 and  $0.6 \text{ g m}^{-3} (\text{km}^{-1} \text{ sr}^{-1})^{-1}$  correspond to the gradient at 2 and  $11 \text{ mg m}^{-3}$ , respectively.

Collocation between MLS and CALIOP measurements was not available during the early period of the CALIOP mission. A critical adjustment in the A-Train formation configuration was made in early 2008. Since then *Aura* MLS and CALIOP footprints are brought within  $\pm 10$  km in the cross-track direction. The A-Train reconfiguration was completed in May 2008, and also improved the MLS-CALIOP temporal separation from  $\sim 7$  to  $< 1$  min. This close alignment among A-Train sensors is critical for cloud studies, because large spatiotemporal variability may exist in upper-tropospheric clouds, especially those from deep convective systems.

In this study we use the version 2 240- and 640-GHz  $T_{\text{cir}}$  that were output as diagnostic products at the end of MLS retrieval process. The variable  $T_{\text{cir}}$  is defined as the difference between the observed and modeled radiances, where the modeled radiance is obtained from the radiative transfer calculation using the best estimated clear-sky atmospheric state (e.g., pressure  $P$  and

temperature  $T$ ) and gas profiles (e.g.,  $\text{H}_2\text{O}$  and  $\text{O}_3$ ). For CALIOP we use the version 3.0 532-nm attenuated total backscatter coefficient  $\beta_{532}$  and IWC data. To match the CALIOP and MLS measurement volumes, we integrate the CALIOP data along the MLS LOS (Fig. 1) to obtain a horizontally integrated IWC [i.e., horizontal ice water path (hIWP)] and integrated 532-nm backscatter  $\gamma_{532}$ , mathematically:

$$\text{hIWP}(h_t) = \int_{-\infty}^{\infty} \text{FOV}(z - h_t) \int_{\text{LOS}} \text{IWC}(s, z) ds dz \quad (1)$$

and

$$\gamma_{532}(h_t) = \int_{-\infty}^{\infty} \text{FOV}(z - h_t) \int_{\text{LOS}} \beta_{532}(s, z) ds dz, \quad (2)$$

where  $\text{FOV}(z)$  is the MLS field of view at tangent height  $z$ , approximately Gaussian, with a frequency-dependent beamwidth. In essence,  $\gamma_{532}$  is a proxy for visible optical depth along MLS LOS after scaled by the lidar ratio (i.e., extinction-to-backscatter ratio). The derived hIWP and  $\gamma_{532}$  are a function of MLS tangent height. At around 15-km tangent height, approximately speaking, hIWP is the IWC integrated over a distance of 200–300 km.

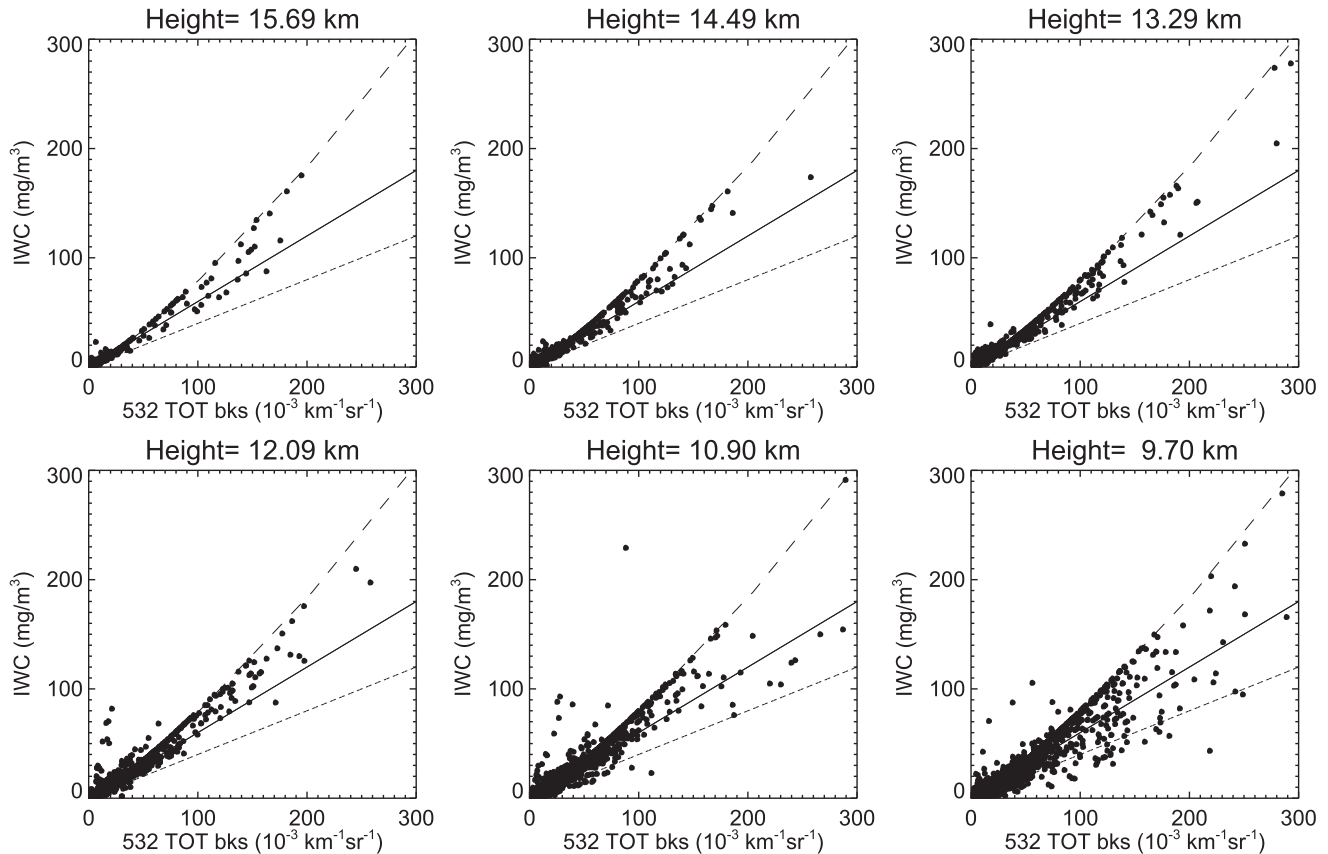


FIG. 2. The IWC-backscatter relationship from CALIOP V3.01 retrievals at selected height levels for 1 Jan 2007. The linear relationships with a slope of 0.4 and  $0.6 \text{ g m}^{-3} (\text{km}^{-1} \text{sr}^{-1})^{-1}$ , and the curve  $\text{IWC} = 0.29\beta_{532}^{1.22}$  are drawn for comparison (see text). The 532-nm backscatter coefficient from CALIOP V3.01 L2\_05 kmCPro dataset is the original attenuated backscatter measurement.

The integrals in Eqs. (1) and (2) neglect the cloud self-extinction effect, which is valid as long as  $\gamma_{532}$  values are small ( $<1$ ) along the MLS LOS. Variability of cloud inhomogeneity along the LOS may increase the noisy nature of comparisons between the matched datasets. However, deep convective cores occur at a much lower frequency than cirrus in the upper troposphere. In this study, we are interested in the cases where  $\gamma_{532} < 1$  and these clouds are assumed to be mostly homogeneous. This volume-matching approach was also used by Wu et al. (2009) in comparing MLS and *CloudSat* cloud ice measurements.

### 3. CALIOP and MLS sensitivities to IWC

Three independent measurements,  $\gamma_{532}$ ,  $T_{\text{cir}}$  (240 GHz), and  $T_{\text{cir}}$  (640 GHz), all sensitive to IWC, are used to evaluate consistency of the microphysics assumptions that lead to their cloud ice retrievals. Here we focus on the 15-km tangent height because CALIOP has relatively good sensitivity to cirrus without much saturation from thick cirrus. Most of the cloud IWC values at 15 km are small ( $<10 \text{ mg m}^{-3}$ ), and therefore it is reasonable to

assume these variables ( $\beta_{532}$ ,  $\gamma_{532}$ ,  $T_{\text{cir}}$ , IWC, and hIWP) are linearly related to each other. Assuming for small cloud perturbations, these variables are basically the first term of the Taylor expansion of radiative transfer equation at a mean atmospheric state. Now, the question is whether these linear relations and underlying microphysics assumptions are consistent.

Without the CALIOP data, it is difficult to determine MLS  $T_{\text{cir}}$  sensitivity to IWC values because of large measurement error in MLS  $T_{\text{cir}}$ . As long as the  $T_{\text{cir}}$  error is random, we can extract its sensitivity to cloud ice from a large volume of MLS  $T_{\text{cir}}$  data by sorting them with respect to the collocated CALIOP  $\gamma_{532}$  measurements. As shown in the density distributions in Fig. 3, a weak linear correlation between  $T_{\text{cir}}$  and  $\gamma_{532}$  at  $\gamma_{532} < 1$  emerges from the statistics of the collocated A-Train data, showing a statistically significant slope with the relative error better than 50% (as seen later in the fitted results). This relationship is expected for thermal emission of cloud ice at microwave frequencies, like those from atmospheric gases. Wu and Jiang (2004) and Wu et al. (2005) studied the sensitivity of MLS 203-GHz limb radiances to cloud ice in the upper troposphere, and

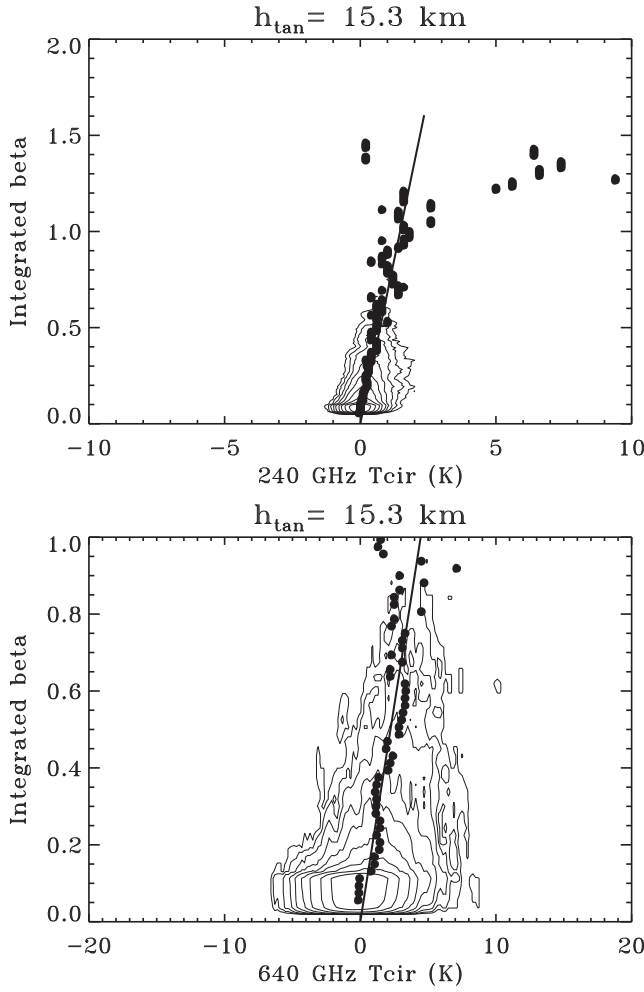


FIG. 3. Observed relations between CALIOP  $\gamma_{532}$  along MLS limb path and MLS 15-km  $\Delta T_{\text{cir}}$  at (top) 240 and (bottom) 640 GHz. A total of approximately 12 000 collocated CALIOP and MLS samples are used. The path integration of CALIOP backscatter along MLS LOS takes into account MLS field-of-view (FOV) effects. Contours are in a logarithmic scale, depicting the density distribution of the measurements. The black dots represent the most probable  $\Delta T_{\text{cir}}$  value at each  $\gamma_{532}$  bin. The line is a fit through these values, producing Eqs. (5) and (6). A bias in  $\Delta T_{\text{cir}}$  ( $-2.2$  K for 240 GHz and  $-1.5$  K for 640 GHz), because of modeled error in the clear-sky radiance, has been removed.

concluded that for small IWCs cloud ice emission can become a dominant process in the  $T_{\text{cir}}$ -hIWP relationship in the tropopause region.

As IWC increases, so does the number of large-size ice particles. As a result, scattering becomes more important in the  $T_{\text{cir}}$ -hIWP relationships. This transition is evident in Fig. 3, where the  $T_{\text{cir}}$ - $\gamma_{532}$  correlation switches drastically from the steep slope to a shallower slope at  $\gamma_{532} > 1$  and  $T_{\text{cir}}(240 \text{ GHz}) > 2$  K. The transition from emission to scattering-dominant  $T_{\text{cir}}$ - $\gamma_{532}$  relationship is frequency dependent because the scattering efficiency varies with frequency. In the MLS version 2.2 (V2.2)

algorithm, only values with  $T_{\text{cir}}(240 \text{ GHz}) > \sim 1.5$  K ( $3\sigma$ ) is considered as useful for cloud ice retrievals, whereas  $T_{\text{cir}}(240 \text{ GHz}) < 1.5$  K or  $T_{\text{cir}}(640 \text{ GHz}) < 5$  K are under the noise floor and classified as “clear sky.” Thanks to CALIOP, now these cloud ice signals can be studied even within the MLS 240- and 640-GHz clear-sky radiances. The MLS measurements are generally consistent with the analysis in Eriksson et al. (2011), who studied the emission-scattering ratio for 348-GHz limb radiances at 14-km tangent height and found that the emission of cloud ice contributes  $\sim 20\%$  to the total  $T_{\text{cir}}$  at large  $T_{\text{cir}}$  values. Their study also suggested an increased ( $\sim 50\%$ ) contribution from ice emission as the 348-GHz limb  $T_{\text{cir}}$  decreases.

In the case where cloud ice is small ( $< 10 \text{ mg m}^{-3}$ ), the  $T_{\text{cir}}$  dependence on IWC can be modeled relatively well from the measured ice dielectric properties at the microwave frequencies. There exists a simple linear  $T_{\text{cir}}$ -IWC relationship in this case. As described in Wu and Jiang (2004), the radiative transfer calculation can be greatly simplified for small ( $< 10 \mu\text{m}$ ) ice crystal and small ( $< 10 \text{ mg m}^{-3}$ ) IWC value situations, under which cloud scattering is negligible. As a result,  $T_{\text{cir}}$  is directly proportional to IWC, similar to the radiance from clear-sky gas emissions. Furthermore, as shown below, we may have an analytical form to characterize the  $T_{\text{cir}}$ -IWC relationship.

Consider the microwave radiances in form of  $T_b = T_0[1 - e^{-(\tau_0 + \Delta\tau_{\text{cir}})}]$ , where  $T_0$  is the ambient air temperature,  $\tau_0$  is the gaseous optical depth along MLS LOS, and  $\Delta\tau_{\text{cir}}$  is the cloud-induced optical depth. For  $\Delta\tau_{\text{cir}} \ll 1$ ,  $T_b$  may be approximately written as

$$T_b \approx T_0[1 - e^{-\tau_0}(1 - \Delta\tau_{\text{cir}})] = T_0[1 - e^{-\tau_0}] + T_0 e^{-\tau_0} \Delta\tau_{\text{cir}},$$

where  $T_{b0} \equiv T_0[1 - e^{-\tau_0}]$  is clear-sky background radiance, and  $T_{\text{cir}}$  is defined by

$$T_{\text{cir}} \equiv T_b - T_{b0} \approx T_0 e^{-\tau_0} \Delta\tau_{\text{cir}}.$$

If scattering is neglected, the cloud-induced optical depth becomes

$$\Delta\tau_{\text{cir}} = 2.1 \frac{\text{IWC}}{\lambda} \frac{3\epsilon''}{(\epsilon' + 2)^2 + \epsilon''^2} \Delta s,$$

where hIWP = IWC  $\cdot$   $\Delta s$  as a simplified form of Eq. (1). Here,  $\Delta s$  is MLS pathlength in kilometers, IWC is in milligrams per cubic meter ( $\text{mg m}^{-3}$ , throughout this paper), wavelength  $\lambda$  is in centimeters, and  $(\epsilon', -\epsilon'')$  are real and imaginary part of ice dielectric constant (Jiang and Wu 2004). At approximately 15 km, within 5% of their variability, the typical values of  $T_{b0}$  in the tropics



are around 40 and around 135 K for MLS 240- and 640-GHz radiances, respectively, and  $T_0 \sim 200$  K is used for the ambient clear-sky air temperature near the tropopause. With these approximate values, we have

$$T_{\text{cir}}(240 \text{ GHz}) \approx 2.4 \times \text{IWC} \times \Delta s \quad (3)$$

and

$$T_{\text{cir}}(640 \text{ GHz}) \approx 7.8 \times \text{IWC} \times \Delta s, \quad (4)$$

which yield a 240- to 640-GHz  $T_{\text{cir}}$  ratio of 1:3.3. Since the ice permittivity is known 10% accuracy at microwave frequencies (e.g., Jiang and Wu 2004), the coefficients in Eqs. (3) and (4) are relatively robust, compared to other error sources as discussed below.

The microwave cloud properties in Eqs. (3) and (4) are further used to verify or constrain CALIOP IWC retrieval in Fig. 2. From Fig. 3 we have

$$\frac{\gamma_{532}}{T_{\text{cir}}} \approx 0.77 (0.31) \quad \text{for } 240 \text{ GHz} \quad (5)$$

and

$$\frac{\gamma_{532}}{T_{\text{cir}}} \approx 0.23 (0.12) \quad \text{for } 640 \text{ GHz}. \quad (6)$$

The number in parentheses is the standard deviation of the fitted slope. Note that the 240- to 640-GHz  $T_{\text{cir}}$  ratio from Eqs. (5) and (6) is 1:3.4, close to the analytical value of 1:3.3 from the simple model in Eqs. (3) and (4). Substituting Eqs. (3) and (4) into Eqs. (5) and (6) and taking the fact that  $\gamma_{532} = \beta_{532} \Delta s$ , we have

$$\frac{\text{IWC}}{\beta_{532}} \approx 0.61 (0.20) \quad \text{for } 240 \text{ GHz} \quad (7)$$

and

$$\frac{\text{IWC}}{\beta_{532}} \approx 0.57 (0.14) \quad \text{for } 640 \text{ GHz}. \quad (8)$$

Both MLS channels suggest that the lidar IWC– $\beta_{532}$  coefficient should be approximately 0.6 to satisfy the microwave cloud properties expected from purely ice thermal emission. By averaging the coefficients in Eqs. (7) and (8), we obtain an empirical coefficients for the IWC– $\beta_{532}$  relationship,

$$\text{IWC} = (0.58 \pm 0.11) \beta_{532}, \quad (9)$$

where  $\beta_{532}$  is in inverse kilometers per steradian. As shown in Fig. 2, the slope of 0.58 from Eq. (9) agrees

quite well with the majority of CALIOP V3.0 retrievals at  $\text{IWC} < 100 \text{ mg m}^{-3}$ . The agreement increases fidelity of the microphysics assumptions used by the CALIOP scattering and MLS emission techniques for cloud ice retrievals at approximately 15 km.

#### 4. Discussion

The linear IWC– $\beta_{532}$  relationship [Eq. (9)] derived in this study provides an independent evaluation on the CALIOP cloud ice retrieval in the V3.01 L2\_05 kmCPro dataset. The key microphysics constraint in Eq. (9), however, is not scattering properties of ice crystals. Rather, it is based on the ice permittivity properties at 240 and 640 GHz, which determine the thermal emission of cloud ice and the  $T_{\text{cir}}$  sensitivity to IWC at these frequencies. In the case where IWC and ice particle sizes are small, scattering can be neglected at the MLS frequencies, which leads to the linear proportionality of  $T_{\text{cir}}$  to IWC. Since ice permittivity is well known at microwave frequencies (with 10% uncertainty; e.g., Jiang and Wu 2004), the uncertainty of Eq. (9) is dominated by the errors in the observed  $T_{\text{cir}}$ – $\gamma_{532}$  relationships in Eqs. (5) and (6).

The agreement between  $\text{IWC} = 0.58 \beta_{532}$  and CALIOP V3.0 retrieval appears to hold well for most altitudes. The increased scatters at lower altitudes in the CALIOP V3.0 data are likely induced by the noise in the extinction retrieval that is expected to increase at lower altitudes in the presence of more high clouds. There is a subtle difference in the IWC– $\beta_{532}$  relation between Eq. (9) and CALIOP V3.0 at  $\text{IWC} < \sim 2 \text{ mg m}^{-3}$ , showing a slope of 0.4, or approximately 50% smaller than 0.58, which would yield a low bias in the CALIOP V3.0 IWC retrieval compared to Eq. (9). At  $\text{IWC} > 11 \text{ mg m}^{-3}$ , the slope in CALIOP V3.0 is greater than 0.6, which would produce a larger IWC retrieval than Eq. (9). The CALIOP V3.0 IWC retrieval employs several assumptions (e.g., lidar ratio, multiple-scattering factor, and extinction-to-IWC coefficient), among which the IWC– $\beta_{532}$  relation was derived from the data with large scatters (Heymsfield et al. 2005). The agreement between CALIOP V3.0 and Eq. (9) suggests that the uncertainty associated with the CALIOP IWC retrieval should be  $< 50\%$ .

To further validate the CALIOP V3.0 and Eq. (9) IWC retrievals, we compare them with *CloudSat* data and statistics of in situ measurements in terms of normalized probability density function (PDF). As described in Wu et al. (2009), the normalized PDF is able to characterize measurement noise, bias, and sensitivity range (sensor noise and saturation), without requiring collocation as long as they have the same ensemble sampling.

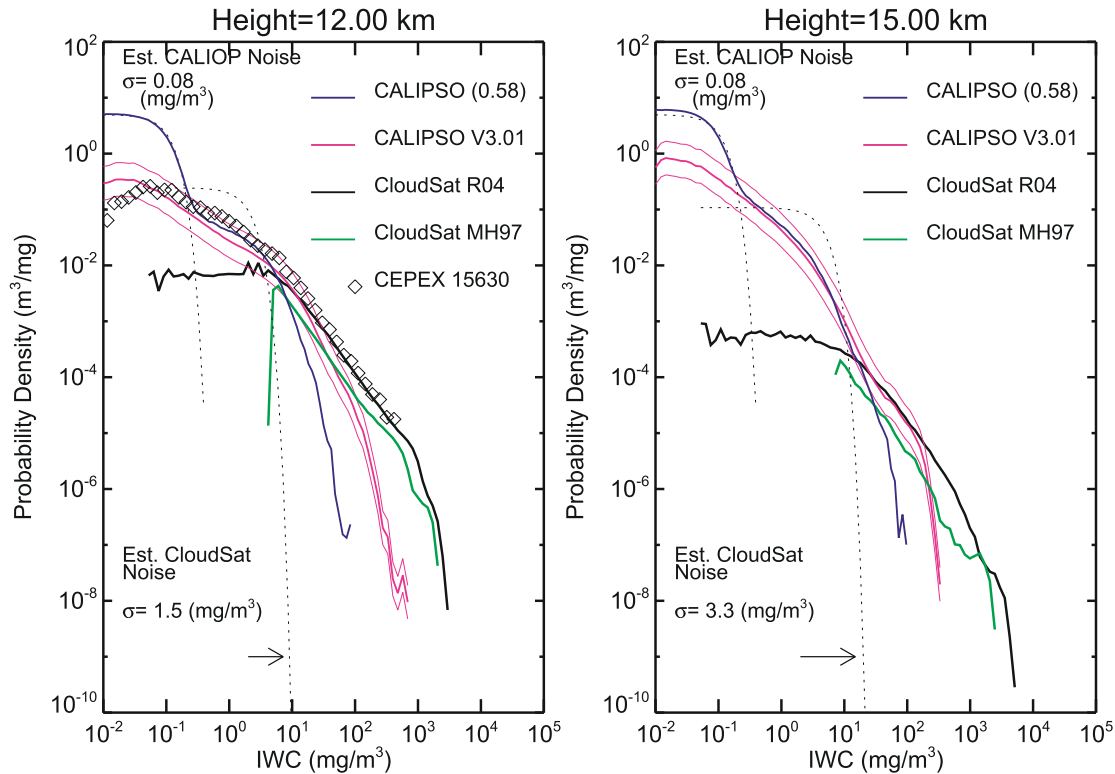


FIG. 4. Normalized PDFs of CEPEX, CALIOP, and *CloudSat* IWC at (left) 12 and (right) 15 km for July 2006 in a tropical band 2.5°S–2.5°N. *CloudSat* IWC retrievals from R04 (Austin et al. 2009) and from the MH97 size distribution are included for comparisons. The CALIOP V3.0 IWC PDF is bound by a scaling factor of 2 to guide comparisons. The CALIOP IWC measurement noise is estimated from the attenuated backscatter assuming the  $IWC = 0.58\beta_{532}$  relation, and the normalized Gaussian noise distribution is the dotted line. Similarly, the *CloudSat* IWC noise is estimated from the reflectivity measurements  $Z_e$  using the calculated  $Z_e$ –IWC relation. The estimated CALIOP and *CloudSat* IWC precisions are 0.08 at 1.5  $\text{mg m}^{-3}$  at 12 km, and 0.08 and 3.3  $\text{mg m}^{-3}$  at 15 km, using the method in Hogan et al. (2006). The PDF from the MH97 retrieval is cut off at  $3\sigma$  to remove the portion below the *CloudSat* reflectivity noise (−27 dBZ).

Here we focus on the CALIOP IWC statistics for latitudes of 2.5°S–2.5°N in July 2006, with comparisons against in situ measurements obtained during Central Equatorial Pacific Experiment (CEPEX) at 12 km, and with *CloudSat* at 12 and 15 km. Two *CloudSat* IWC retrievals are included: one from the standard IWC product in the R04 release (Austin et al. 2009) and the other from the retrieval assuming the MH97 size distribution (Eriksson et al. 2008; Rydberg et al. 2009).

As seen in Fig. 4, there is good agreement between CALIOP and *CloudSat* R04 IWC PDFs at 12 and 15 km in the overlapped sensitivity ranges. The two sensors overlap for IWC values of 5–20  $\text{mg m}^{-3}$  at 12 km and 30–200  $\text{mg m}^{-3}$  at 15 km, respectively. The overlapped sensitivity requires the sensitivity from both instruments must be greater than their measurement noise but not saturated. Figure 4 shows that CALIOP PDFs drop off sharply at large IWC values, as expected for the increased attenuation by the dense cloud at a higher altitude. The attenuation correction, as implemented in

the CALIOP V3.0 algorithm, can mitigate the problem to the extent where clouds are moderately thick. Thus, we should focus more on the comparison at small IWC values, where clouds are relatively thin and above the *CloudSat* detector noise (−31 dBZ). The *CloudSat* cloud ice PDF agrees with CALIOP V3.0 at 5–20  $\text{mg m}^{-3}$  for 12 km and at 30–200  $\text{mg m}^{-3}$  for 15 km. The reduced CALIOP sensitivity overlap with *CloudSat* at 12 km is, on one hand, a manifestation of increased attenuation from the clouds above. On the other hand, it is limited by the lower limit in *CloudSat* cloud detection. The *CloudSat* IWC noise produces a white PDF below its detection limit, corresponding to approximately 5  $\text{mg m}^{-3}$  at 12 km and approximately 30  $\text{mg m}^{-3}$  at 15 km, respectively.

Despite the reduced sensitivity overlap with *CloudSat* at 12 km, it is encouraging to observe the agreement between *CloudSat* and CALIOP cloud ice probability at the 5–20  $\text{mg m}^{-3}$  range. More importantly, cloud ice statistics are extended for most cirrus with IWC between approximately 5  $\text{mg m}^{-3}$  ( $3\sigma$  *CloudSat* noise) and

approximately  $0.2 \text{ mg m}^{-3}$  ( $3\sigma$  CALIOP noise). The CALIOP V3.0 PDF shows a slightly low bias against CEPEX PDF at IWC values between 0.5 and  $5 \text{ mg m}^{-3}$ . A plausible cause could be the cloudy-sky bias from the sampling during the CEPEX campaign. Another possibility, as aforementioned, is because of a smaller slope in the IWC- $\beta_{532}$  conversion at  $\text{IWC} < \sim 2 \text{ mg m}^{-3}$ , which would lower the probability at IWC near this value. On the other hand, the retrieval using Eq. (9) would produce a probability closer to the CEPEX statistics.

The 12 km is a critical altitude for *CloudSat* IWC validation against in situ measurements. As pointed out in Wu et al. (2009), the *CloudSat* R04 IWC statistics agree well with the CEPEX statistics over a broad range of IWC ( $5\text{--}1000 \text{ mg m}^{-3}$ ). In a comparison of MLS V2 and *CloudSat* R04 IWC, Wu et al. (2009) found that MLS mean IWC is lower by a factor of approximately 5 against *CloudSat* at 147 and 100 hPa. The microphysics assumption was thought to be the key cause of this difference. To further evaluate the impacts of the microphysics assumption on *CloudSat* IWC retrievals, we compare the IWC retrievals using the MH97 parameterization, as used in the MLS retrieval, to the R04 product (assuming a gamma size distribution). As shown in Fig. 4, the MH97 IWC retrieval lowers the PDF by a factor of approximately 2 and approximately 4 at 12 and 15 km, respectively. This is generally consistent with the study by Eriksson et al. (2008), showing that the mean IWC is lower by 1.3 at 12.5 km and 2.4 at 15.5 km. All these cross-satellite evaluations support the speculation that the lower bias in MLS IWC, which increases with height, was because of the MH97 size distribution assumption (Wu et al. 2009). On the other hand, the agreement of cloud ice measurements among *CloudSat*, CALIOP V3.0, and Eq. (9) at 15 km suggests that the gamma size distribution, as used by *CloudSat*, appears to be more realistic for upper-tropospheric ice clouds.

## 5. Conclusions

Analyzing a large ensemble of collocated A-Train CALIOP backscatter  $\beta_{532}$  and MLS cloud-induced radiance  $T_{\text{cir}}$  measurements, we obtained the empirical relationships among  $\beta_{532}$  and MLS 240- and 640-GHz  $T_{\text{cir}}$ . These linear relations are statistically significant and lead to an empirical extinction-to-IWC conversion in Eq. (9) for small IWC values, that is,  $\text{IWC}/\beta_{532} = 0.58 \pm 0.11$ . The key microphysics assumption in Eq. (9) is the ice permittivity property at microwave frequencies, which is known to a good accuracy. Because of the noisy MLS measurements in small  $T_{\text{cir}}$  values, the uncertainty in Eq. (9) is dominated by the observed  $T_{\text{cir}}\text{--}\gamma_{532}$

relationships. Since MLS  $T_{\text{cir}}$  is directly proportional to IWC and independent of the shape of particle size distribution for small IWC values, the resulting IWC- $\beta_{532}$  relation provides additional constraint on the CALIOP cloud ice retrieval.

The empirical IWC- $\beta_{532}$  relation in Eq. (9) agrees well with the extinction-to-IWC conversion used in the CALIOP V3.0 retrieval. This agreement improves fidelity of the scattering-based CALIOP and emission-based MLS IWC retrievals, as a result of the consistent microphysics in explaining the observed  $T_{\text{cir}}\text{--}\gamma_{532}$  correlation. Furthermore, the agreement between CALIOP and *CloudSat* IWC PDFs suggests that the MH97 parameterization be the primary cause of the underestimation of MLS cloud ice at 15 km and the altitudes above.

Finally, we demonstrate in this study that multisensor analyses can be used to constrain the microphysics assumptions used in satellite cloud ice retrievals. Through self-consistency evaluation on the collocated A-Train measurements, we are able to cross evaluate the cloud ice measurements at two extreme wavelengths: millimeter- and submillimeter-wave versus visible. Both emission-based microwave radiometry and scattering-based lidar backscattering reach a statistically consistent ( $<50\%$ ) IWC- $\beta_{532}$  relationship on different assumptions about ice cloud microphysics in the upper troposphere. This agreement raises fidelity on these assumptions and sheds a new light on the issues unsolved by earlier cloud validation efforts. Similar studies can be applied to other A-Train sensor pairs or closely sampled datasets from formation flights.

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